

## Folding during soft-sediment deformation

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### Abstract

The detailed analysis of folding in rocks was in part pioneered by John Ramsay and resulted in a range of techniques and criteria to define folds. Although folding of unlithified or 'soft-sediments' is typically assumed to produce similar geometries to those in 'hard rocks', there has to date been little detailed analysis of such folds. The aim of this paper is therefore to investigate folds developed during soft-sediment deformation (SSD) by applying techniques established for the analysis of tectonic folds during hard rock deformation (HRD). We use the late-Pleistocene Lisan Formation exposed around the Dead Sea as our case study, as the laminated lake sediments record intricacies of fold detail generated during seismically-triggered slumping of mass transport deposits (MTD's) towards the depocentre of the basin. While it is frequently assumed that folds created during SSD are chaotic and form disharmonic structures, we provide analyses that show harmonic fold trains may form during slumping, although larger upright folds cannot be traced for significant distances and are more typically disharmonic. Our analysis also reveals a range of fold styles, with more competent detrital-rich layers displaying buckles (Class 1B) as well as upright Class 1A folds marked by thickened limbs. Class 1A buckle folds are generally considered to be created by flattening that overprints folds with original Class 1B geometry. As thickened fold limbs are truncated by overlying erosive surfaces, the vertical flattening is considered to have occurred during the slump event. Different fold shapes may partially reflect variable flattening, depending on the original orientation of upright or recumbent folds, together with continued downslope-directed simple shear deformation that modifies the fold geometry. Analysis of fold wavelength, amplitude and bed thickness allows us to plot strain contour maps, and indicates that beds defining slump folds display viscosity contrasts in the range of 50 – 250, which are similar to values estimated from folds created during HRD in metamorphic rocks. A range of re-fold patterns, similar to those established by John Ramsay in metamorphic rocks, are observed within slumps, and are truncated by the overlying sediments indicating that they formed during a single progressive slump event rather than distinct 'episodes' of superimposed deformation. This study confirms that techniques developed for the analysis of folds created during HRD are equally applicable to those formed during SSD, and that resulting folds are generally indistinguishable from one another. Extreme caution should therefore be exercised when interpreting the origin of folds in the rock record where the palaeogeographic and tectonic contexts become increasingly uncertain thereby leading to potential misidentification of folds created during SSD.

**Keywords:** folding, slumps, mass transport deposit, Dead Sea

43 It is generally assumed that slump folds generated during soft-sediment deformation (SSD)  
44 within mass transport deposits (MTDs) bear a close resemblance to tectonic folds created  
45 during hard rock deformation (HRD) (see discussions in Elliot & Williams 1988; Maltman  
46 1994; Alsop et al. 2019a). However, there has in reality been very little systematic analysis of  
47 fold styles created during soft-sediment deformation, with some notable exceptions including  
48 Woodcock (1976a, b, 1979), Farrell & Eaton (1987), Waldron & Gagnon (2011) and Ortner  
49 & Kilian (2016). Hudleston (1986, p.238) noted that “shape alone cannot be used to  
50 distinguish soft-sediment folds from tectonic folds” (see also Woodcock 1976b). Comparison  
51 of curvilinear fold patterns in metamorphic rocks and soft-sediments also reveals no  
52 significant geometric difference between aspects of curvilinear fold geometries (e.g. Alsop et  
53 al. 2007).

54 While folding within metamorphic rocks is generally created by deformation  
55 associated with recrystallization, folding within unlithified or ‘soft sediments’ is typically  
56 achieved via independent particulate flow, where individual grains move relative to one  
57 another to create a range of structures (e.g. Knipe 1986). The exact nature of these structures  
58 is dependent on the ratio of pore fluid pressure to cohesive strength of the sediment (e.g.  
59 Knipe 1986; Ortner 2007), with folds and shears being created where fluid pressure is less  
60 than cohesive strength leading to hydroplastic deformation. The examination of fold  
61 geometries is important as Waldron & Gagnon (2011) realised that fold styles in different  
62 sand and mud layers may be used to distinguish the degree of lithification during the fold  
63 process, and thereby identify ‘soft-sediment’ folds. Despite the obvious differences in the  
64 mechanisms of folding within HRD and SSD noted above, the geometry of the resulting  
65 structures are so similar to one another that criteria to distinguish between these different  
66 types of folds in ancient settings should be considered (see Alsop et al. 2019a for a review).

67 The broad aim of this contribution is to document and investigate fold styles created  
68 during SSD by applying classical techniques of structural geology such as dip-isogon analysis  
69 and examination of fold interference patterns established by John Ramsay among others (e.g.  
70 Ramsay 1967, Ramsay & Huber 1987). Van der Pluijm & Marshak (2004, p.25) discuss  
71 slump folds and note that “folds in one layer are of a different size and orientation than the  
72 structures in adjacent layers” suggesting a largely disharmonic style. We therefore examine  
73 trains of soft-sediment folds to determine whether they are indeed disharmonic using  
74 techniques of fold spacing compared to distance they can be traced along their axial surface  
75 (e.g. Twiss & Moores 2007, p.290). We also analyse folded layer thickness, amplitude and  
76 wavelength in an attempt to estimate the viscosity contrast between layers during sediment  
77 folding together with % shortening (e.g. Schmalholz & Podladchikov 2001). Our analysis  
78 aims to help answer the following research questions.

- 79 i) Does slumping create harmonic or disharmonic folds?  
80 ii) Are refold patterns in slumps similar to those in metamorphic rocks?  
81 iii) What range of fold styles are created during slumping?  
82 iv) How are slump folds subsequently modified?  
83 v) Are estimates of viscosity contrasts in slumps folds similar to those in metamorphic rocks?

84 We first describe the regional setting of the case study area in the Dead Sea Basin, before  
85 analysing folds created during SSD linked to downslope-directed slumping of MTDs towards  
86 the basin.

87

## 88 **Regional setting**

89 The Dead Sea Basin is a pull-apart structure on the Dead Sea Fault, which is marked  
90 by two major parallel fault strands that generate numerous earthquakes (Fig. 1a, b) (e.g.  
91 Marco et al. 1996, 2003; Ken-Tor et al. 2001; Migowski et al. 2004; Begin et al. 2005). This  
92 transform is thought to have been active from the Miocene to recent, including during  
93 deposition of the Late Pleistocene (70-15 Ka) Lisan Fm. that forms the focus of the present  
94 study (e.g. Bartov et al. 1980; Garfunkel 1981; Haase-Schramm et al. 2004). The Lisan Fm.  
95 comprises mm-scale aragonite laminae that were precipitated from the upper waters of Lake  
96 Lisan during the summer, together with more detrital-rich layers washed into the lake during  
97 flood events (Begin et al. 1974) (Fig. 2a, b). The detrital units have been sampled < 1 km NE  
98 of the Peratzim site by Haliva-Cohen et al. (2012) (see Fig. 1b), and compositionally consist  
99 of quartz and calcite grains with minor feldspar, and clays (illite-smectite). Detrital laminae  
100 within the varved aragonite-rich Lisan Fm. display grain sizes of ~8-10  $\mu\text{m}$  (silt), while the  
101 thicker (> 10 cm) detrital-rich units are generally coarser grained (60 – 70  $\mu\text{m}$ ) and can be  
102 classified as very fine sands (Haliva-Cohen et al. 2012). Although deposited on slopes of <1°,  
103 the Lisan Fm. contains numerous intraformational fold and thrust horizons that are capped by  
104 undeformed beds and are considered to be seismically triggered (Alsop & Marco 2013; Alsop  
105 et al. 2016; Lu et al. 2017) (Fig. 2a). Some evidence for bedding parallel shear during seismic  
106 events is also preserved by offset clastic dykes (Weinberger et al. 2016).

107 The Lisan Fm. is exposed for ~100 km along the western margin of the Dead Sea  
108 Basin and displays systematic variation in the orientation of slump fold and thrust systems  
109 within it. In the northern portions slumping is directed towards the ESE, in the central portion  
110 around Masada slumping is towards the east, whilst in the southern area around Peratzim  
111 slumping is NE-directed (Alsop & Marco 2012a) (Fig. 1b). Combined with westerly directed  
112 slump folds recorded from the eastern shore of the Dead Sea in Jordan (El-Isa & Mustafa  
113 1986), this suggests a regional pattern of radial slumping towards the depocentre of the Dead  
114 Sea Basin (Alsop & Marco 2012a). In the extreme southern area around Zin, slumping is  
115 towards the south and is interpreted to be influenced by the nearby NE-SW trending  
116 Amazyahu Fault (Weinberger et al. 2017, Alsop et al. 2018; 2019b) (Fig. 1b). This overall  
117 pattern has been subsequently corroborated by Anisotropy of Magnetic Susceptibility (AMS)  
118 fabrics that support a radial pattern of MTD slumping (Weinberger et al. 2017).

119 The Dead Sea Basin is an ideal place to study structures associated with SSD as the  
120 relatively cohesive muds and precipitated aragonitic layers define an intricate varve-like  
121 stratigraphy that define a range of detailed structures that may not survive elsewhere (Fig. 2a,  
122 b). The MTD horizons and reworked zones of sediment we are particularly interested in for  
123 this case study are best exposed around Masada [N31°1847.0 E35°2228.3], Peratzim  
124 [N31°0449.6 E35°2104.2], and Zin [N31.00615, E35.26342] in the central and southern Dead

125 Sea areas respectively (Fig. 1b) (see Weinberger et al. 2017). All of these sites are positioned  
126 ~1 km east of the Dead Sea western border fault zone, with Cenomanian-Senonian carbonates  
127 preserved further to the west in the footwall to this fault. For most of the time between 70 and  
128 28ka, Lake Lisan in these areas had a maximum depth of 100 m or less, apart from a brief  
129 period from 26-24 ka when water was up to 200 m deep (Bartov et al. 2003).

130

### 131 **Analysis of slump fold and thrust displacement**

132 Six individual slump sheets that form MTDs have been recognised in the Peratzim area  
133 (Alsop et al. 2016). Although some variability exists, slump fold hinges are typically NW-SE  
134 trending and verge and face towards the NE, with transport directed towards the depocentre  
135 of the basin (Alsop et al. 2016) (Fig. 2c-h). Fold axial planes dip variably towards the SW,  
136 while downslope-verging fore thrusts dip towards the SW (Fig. 2c-h) (Alsop et al. 2017a). At  
137 the Masada locality, fold hinges are typically NNW-SSE trending with gently WSW dipping  
138 axial planes and overall ENE-directed vergence with transport towards the depocentre of the  
139 basin (Alsop & Marco, 2012a) (Fig. 3a, b). At the Zin locality, fold hinges are E-W trending  
140 and typically verge and face towards the south (Weinberger et al. 2017; Alsop et al. 2018).  
141 As with previous studies, care has been taken that cuttings along wadi walls are developed at  
142 high angles to fold hinges, thereby providing transport-parallel (or hinge-normal) profile  
143 sections to undertake structural analysis (e.g. Alsop et al. 2017b).

144 Each MTD horizon is overlain by undeformed sediments that were deposited after  
145 each slope failure event (e.g. Fig. 2a). The tops of MTDs, and the folds they contain, may be  
146 truncated by the overlying sedimentary ‘cap’ that was deposited out of suspension following  
147 the failure event (Alsop & Marco 2012b, Alsop et al. 2016, 2019a, b). This cap, which may  
148 be up to 20 cm thick, blankets and infills the underlying slump topography, and demonstrates  
149 that the structures formed at or close to the sediment surface. The studied sections form in the  
150 upper part of the Lisan Fm. and have never had significant overburden (< 10 m) above them.  
151 As such, the structures created during slumping have not undergone significant later  
152 modification, and preserve pristine geometries unmodified by later compaction (Alsop et al.  
153 2017a, b, 2019a). In addition, the aragonite- and detrital-rich alternations within the Lisan  
154 Fm. form a bilaminate sequence i.e. “comprising only two different types of layers which  
155 alternate with each other” (Price & Cosgrove 1990, p.307). Such bilaminates simplify the  
156 analysis of folding, although they need not be regularly spaced (equal thickness) and may  
157 form either multilayer packages of folding (where folding is closely spaced) or single layer  
158 buckles where competent beds are separated from one another by a weaker ‘host’.

159

### 160 **Harmonic and disharmonic slump folding**

161 A harmonic fold is simply defined as being continuous along its axial surface for ‘many  
162 multiples of the half wavelength’ whereas a disharmonic fold ‘dies out within a couple of half  
163 wavelengths’ (Twiss & Moores 2007, p.289; see also Fossen 2016 p.259). The half  
164 wavelength may be approximated by the spacing (S) of adjacent axial surfaces, while the

165 continuity of the axial surface may be directly measured (D), to produce a ratio (H) which is  
166 equivalent to D/S (see Twiss & Moores 2007, p.290) (Fig. 3a).

167 Upright folds within the Lisan Fm. are frequently disharmonic, as they are positioned  
168 above basal detachments and are overlain by sedimentary caps meaning that they cannot be  
169 traced for any significant distance (D) along their axial surfaces (e.g. Fig. 2c, e, g). However,  
170 we have also analysed trains of recumbent slump folds (Fig. 3a, e, i). In general the maximum  
171 distance (D) that individual decimetric-scale folds may be traced along their axial planes  
172 increases as the maximum spacing (S) between adjacent axial planes becomes greater (Fig.  
173 3a, c). In addition, the ratio H (D/S) ranges from ~18 to ~6, and is shown to reduce in a non-  
174 linear way as S increases (Fig. 3d). Analysis of individual fold systems reveals broadly  
175 similar patterns of D, S, and resulting H (Fig. 3e-l). Individual fold packages may form en  
176 echelon fold trains that migrate up through the folded sequence where they cross thicker  
177 aragonite-rich units (Fig. 3a). These fold packages do not directly interact with one another  
178 and do not therefore create re-fold patterns (see section 5) below). Overall, these recumbent  
179 folds may be described as broadly harmonic as  $D \gg S$ .

180

### 181 **Refold patterns within individual slump horizons**

182 Working on polyphase deformed metamorphic rocks in the Scottish Highlands, Ramsay  
183 (1962, 1967 p.518) developed a classification scheme of refolded folds, based on the relative  
184 orientations of different sets of fold hinges and axial surfaces that were superimposed on one  
185 another and considered to develop sequentially during orogenesis. Ramsay's (1962, 1967)  
186 classic Loch Monar dam outcrops comprise superb water-washed exposures in psammites of  
187 the Precambrian Moine Supergroup, and permit intricate details of folding to be examined  
188 (e.g. Fig. 4a). Ramsay (1967 p.521) recognised Type 1 or 'dome and basin' re-fold patterns  
189 created by superimposition of fold hinges and axial planes at high angles (~90°) to one  
190 another (Fig. 4b, c). Type 2 crescent or 'angel-wing/ mushroom' styles of refolding formed  
191 by superimposition of fold hinges and axial planes at moderate to high angles to one another  
192 were also observed (Ramsay 1967, p.525), as were Classic Type 3 'hook' refolds that form  
193 where superimposed fold hinges are coaxial, while axial surfaces are at high angles to one  
194 another (Ramsay 1967, p.530) (Fig. 4d). This general classification scheme has subsequently  
195 been applied across a huge range of metamorphic rocks, scales and settings.

196 Within the Peratzim area of the Dead Sea case study (Fig. 1b), Type 1 or 'dome and  
197 basin' re-fold patterns are developed in wadi sections within individual MTDs (Fig. 4e, f).  
198 Type 2 crescent or 'angel-wing/ mushroom' styles of refolding are also observed (Fig. 4g) as  
199 are classic Type 3 'hook' refolds that are particularly common (Fig. 4h, i, j). Refold patterns  
200 are truncated by the overlying sedimentary cap that was deposited from suspension following  
201 the slump event (Alsop & Marco 2012b) (e.g. Fig. 4j). This indicates that refolding was  
202 completed during the slump event and was not a much later re-working created via  
203 subsequent loading or later slumping. Thus, a re-working or modification of folds to create  
204 classic re-fold patterns occurs during a single, progressive slump event.

205

## 206 **Dip-isogon analysis around slump folds**

207 The dip-isogon method is a well-established technique of fold classification in lithified rocks  
208 developed by John Ramsay (e.g. Ramsay 1967, p.363). In this method, dip isogons join  
209 points of equal dip on adjacent folded surfaces within the fold profile,  $t_0$  is layer thickness  
210 measured along the axial surface, while  $t_\alpha$  is orthogonal layer thickness measured at various  
211 angles ( $\alpha$  - alpha) to the reference plane (Fig. 5a). Graphs normalise thicknesses by using  $t'_\alpha$   
212 (where  $t'_\alpha = t_\alpha / t_0$ ) and plot this value against dip angle ( $\alpha$ ) to create a series of fold classes  
213 (Ramsay 1967, p. 366). Class 1 folds are marked by convergent dip isogons (Fig. 5b), Class 2  
214 folds by parallel dip isogons (Fig. 5c), and Class 3 folds by diverging dip isogons (Fig. 5d)  
215 (e.g. Ramsay 1967, p.365; see Fossen 2016, p.263).

216 We use the dip-isogon method to analyse and compare fold geometries formed in  
217 aragonite-rich and detrital-rich units within the case study at Peratzim (Fig. 1b). The  
218 unlithified sediment enables easy excavation, and reveals that E-W trending asymmetric  
219 buckle folds with gently south-dipping axial planes verge towards the north (Fig. 5 e, f, g).  
220 Our analysis includes data from both the upper and lower limbs of folds, and shows that folds  
221 within aragonite-rich units display gently convergent to parallel dip isogons that typically  
222 define Class 1C folds (Ramsay 1967, p.367; Fossen 2016, p.263) (Fig. 5h, i). Conversely, the  
223 detrital units are marked by a strongly convergent isogon pattern representing Class 1C or 1B  
224 parallel folds consistent with buckling (Fig. 5h, i). In detail, the upper limbs are in part Class  
225 1C (Fig. 5h, i), while the lower overturned limb may exhibit a component of thickening  
226 relative to the hinge area to create Class 1A folds (Fig. 5h, i, 6a, b). However, as folds  
227 become progressively more asymmetric and overturned, the lower limb is consequently  
228 thinned (Fig. 5e). These overturned buckles also display a reduction in the angle of axial  
229 planar dip from  $70^\circ$  to  $24^\circ$ , together with a slight re-orientation of fold hinges from  $092^\circ$   
230 towards  $063^\circ$  and the general slope direction at Peratzim (Fig. 2c-h, 5e,f,g). In general, more  
231 upright folds within detrital-rich units display Class 1A or Class 1B (parallel) buckle fold  
232 styles, while recumbent folds are marked by Class 1C or Class 2 (similar) geometries (Fig.  
233 6a-h). Aragonite-rich units generally display a parallel isogon pattern most consistent with  
234 Class 2 similar folding (Fig. 6c, d, g, h). These relationships collectively indicate that detrital-  
235 layers were locally more competent (see also Alsop et al. 2017a, b), and suggest that folding  
236 involved a component of progressive deformation that resulted in tightening and rotation of  
237 fold hinges towards the flow direction, while axial planes were rotated into the plane of flow.

238 Ramsay (1967 p.432) recognised that multilayers folds may be composed of adjacent  
239 beds defining different classes of folds such as Class 1C in more competent units and Class 3  
240 in the incompetent horizons, although the overall fold represent the sum of these layers and  
241 will closely resemble a Class 2 similar fold. Within individual folds of the Lisan Fm.,  
242 different layers define different classes of fold, with aragonite-rich layers displaying Class 2  
243 (similar) folds, while adjacent detrital-rich beds are marked by Class 1C fold geometries (Fig.  
244 6g, h). Such varying styles of folding allow an individual fold to propagate further along its  
245 axial plane without encountering significant accommodation problems around the fold hinge  
246 (e.g. Price & Cosgrove 1990, p.320).

247

## 248 **Post-buckle modification of slump folds**

249 Buckling is simply defined as “the flexing or folding of a surface or series of parallel surfaces  
250 by a compressive stress directed along that surface or layer” (Price & Cosgrove 1990, p.273).  
251 However, Ramsay (1967, p.434) also recognised that buckle folds may undergo subsequent  
252 deformation and flattening that transforms their geometry from typical Class 1B parallel folds  
253 to Class 1C or Class 2 (similar) folds. A simple technique to determine the amount of post-  
254 buckle flattening was developed by Lisle (1992), and uses the inverse thickness method,  
255 where  $t$  is the orthogonal thickness measured between two tangents for the inner and outer  
256 layer boundaries. The inverse layer thickness ( $1/t$ ) is then plotted for various orientations of  
257 the layer tangent around the fold (see Lisle 1992, p.370) (Fig. 7a, b). This method has the  
258 significant advantage that post-buckle flattening does not have to be aligned parallel to the  
259 axial trace of the buckle fold (Lisle 1992). Although this technique was devised and applied  
260 to folds formed during HRD (e.g. Lisle 1992; Alsop et al. 1998, their fig. 7), the present study  
261 forms its first use on folds created during SSD. Within the study area, elliptical ratios using  
262 the method of Lisle (1992) range between 1.77 (Fig. 7a, b), 1.70 Fig. 7e, f), 1.97 Fig. 7i, j),  
263 suggesting a relative vertical shortening and flattening of ~30%.

264 Within the Lisan Fm., folding of detrital-rich marker layers that are more competent  
265 than surrounding aragonite-rich horizons (e.g. Alsop et al. 2017a, b) results in a parallel  
266 (Class 1B) style of folding (Fig. 5h,i). The flattening of upright Class 1B folds by vertical  
267 shortening results in thicker limbs compared to hinges (Class 1A) (Figs. 6a,b, 7a-l), whereas  
268 potential vertical flattening of recumbent horizontal folds results in limbs becoming thinner  
269 (relative to hinges) and creation of Class 1C or similar (Class 2) folds (Fig. 6c-h) (see Farrell  
270 & Eaton 1987, their figure 11). Examples of both types of folds defined by detrital-rich  
271 marker horizons that are originally considered to have formed Class 1B folds are observed in  
272 the Lisan Fm. (Fig. 6c-h). Clearly, the post-buckle flattening component is independent of the  
273 orientation of the buckle fold, with the orientation of the strain ellipse indicating that  
274 flattening was consistently sub-vertical (Fig. 7a-l). The thickening and thinning patterns  
275 observed around the slump folds are similar to the theoretical models of flattened folds as  
276 discussed by Twiss & Moores (2007, p.374).

277 An elegant development to the model of Lisle (1992) was provided by Srivastava &  
278 Shah (2006) who realised that strain ellipses (and the associated photograph of the fold) may  
279 simply be ‘unstrained’ on drafting packages by restoring the calculated strain ellipse back to  
280 an original circle i.e. Fig. 7a and 7b are transformed to Fig 7c and 7d respectively when  
281 ‘unstrained’ (Fig. 7a-l). The upslope and downslope limbs have slightly different ellipse  
282 shapes and this perhaps reflects the original limb thicknesses that vary as a result of fold  
283 vergence downslope towards the depocentre. The observation that upright Class 1A folds  
284 with thickened fold limbs are developed close to the sediment surface and are truncated by  
285 the overlying sedimentary cap (Fig. 7i-l) demonstrates that flattening formed during the  
286 actual slump event, and is not a later effect created by loading from overlying MTDs.

287

## 288 **Analysis of layer thickness, amplitude and wavelength in slump folds**

289 Layer thickness ( $h$ ), amplitude ( $A$ ) and wavelength ( $\lambda$ ) of single layer folds (rather than multi  
290 layers) may be measured in an attempt to estimate the strain and viscosity contrast between  
291 layers during folding (e.g. Schmalholz & Podladchikov 2001; see also Hudleston & Treagus  
292 2010). The technique involves analysis of single layer folds (i.e. unaffected by neighbouring  
293 competent beds) and is based on results for linear viscous folding rather than power law  
294 viscous folding (Schmalholtz & Podladchikov 2001). Thickness of a layer ( $h$ ) is measured  
295 orthogonal to the folded layer, while amplitude ( $A$ ) is defined as half the distance from the  
296 trough to the crest of upright folds (e.g. Fig. 8a). Wavelength ( $\lambda$ ) is defined as the distance  
297 between two points that occupy a similar position on the fold train (i.e. between adjacent  
298 synform hinges) (e.g. Fig. 8a). Schmalholz & Podladchikov (2001, p. 206) state that  
299 wavelength may also be measured as double the horizontal distance between neighbouring  
300 fold hinges (i.e. double the distance between antiform and synform fold hinges forming a fold  
301 pair). In this method of fold analysis, amplitude / wavelength ( $A/\lambda$ ) is compared with layer  
302 thickness / wavelength ( $h/\lambda$ ) on a strain contour map (Schmalholz & Podladchikov 2001). For  
303 any fold where amplitude, thickness and wavelength can be measured on the profile plane,  
304 estimates of bulk strain (in terms of % shortening) and the layer / matrix viscosity ratio can  
305 be made by reading the position of data directly off the map (e.g. Fig. 8a-r). This study forms  
306 the first use of this methodology in folds created during SSD, and assumes all the layer  
307 shortening is taken up by buckling with no out of plane movement.

308 Within individual fold trains, wavelength of folds ( $\lambda$ ) reduces as amplitude ( $A$ )  
309 increases, so that the  $A/\lambda$  ratio defines a general trend when plotted against  $\lambda$  (Fig. 8a-b, d-e,  
310 g-h, j-k, m-n, p-q). In detail, when compared to wavelength, the  $A/\lambda$  ratio is not a straight line,  
311 with amplitude increasing more slowly than wavelength (Fig. 8e). In general, steeper folds  
312 typically have lower  $A/\lambda$  ratios, as recumbent folds can develop proportionally greater  
313 amplitudes due to being unconstrained and unhindered by the overlying sediment surface  
314 (e.g. compare upright folds in Fig. 8g, h with recumbent folds in Fig. 8p, q). The higher  $A/\lambda$   
315 ratios of recumbent folds may also be a product of increased simple shear, which rotates the  
316 axial plane towards the shear plane (e.g. Fig. 5e,f,g, 8a-c). Overall analysis of individual fold  
317 data sets on the strain contour maps (Schmalholz & Podladchikov 2001) reveals that most  
318 folded layers display viscosity contrasts in a range between 50 and 250, while calculated  
319 layer shortening is generally between 30% and 70%, (Fig. 8c, f, i, l, o, r). Within individual  
320 fold trains, % contraction typically increases as folds become more inclined, or recumbent  
321 (e.g. Fig. 8a-c), with some small-scale recumbent folds recording values in excess of 70%  
322 (e.g. Fig. 8p-r). When plotted on strain contour maps, the overall trends of data from  
323 individual fold trains are slightly oblique to the established lines marking fixed viscosity  
324 contrasts, with some plots suggesting folds with lower % shortening are marked by lower  
325 viscosity contrasts compared to folds with higher % shortening (Fig. 8c). In other more  
326 typical cases, folds with lower % shortening have greater viscosity contrasts compared to  
327 adjacent folds with higher % shortening, resulting in more 'gentle' trends on strain contour  
328 maps (Fig. 8f, i, l, o).

329 When the combined data set is considered, amplitude ( $A$ ) shows a clear correlation  
330 with wavelength ( $\lambda$ ) across a range of scales from mm (Fig. 9a), to cm (Fig. 9b) to larger  
331 folds shown on log-log plots (Fig. 9c). Similarly, layer thickness ( $h$ ) and fold wavelength ( $\lambda$ )



332 increase in tandem and correlate across a range of scales from mm to m (Fig. 9d-f), as do  
333 layer thickness (h) and fold amplitude (A) (Fig. 9g,h). These general correlations produce a  
334 cluster of points on the overall strain contour plot, suggesting viscosity contrasts in the range  
335 of 50-250, and % shortening between 30 and 70% (Fig. 9i).

336

## 337 Discussion

### 338 *Does slumping create harmonic or disharmonic folds?*

339 Our results demonstrate that the distance (D) fold trains can be traced along their axial  
340 surfaces, compared to the spacing (S) with the neighbouring axial surface show distinct and  
341 consistent relationships (Fig. 3c, g, k). The ratio H (D/S) varies between ~6 and 18, and  
342 generally increases as the spacing (S) reduces (Fig. 3d, h, i). These relationships collectively  
343 demonstrate that overturned and recumbent slump folds that can be traced for < 1 m display  
344 harmonic fold relationships (Figs. 3, 10). Price & Cosgrove (1990 p.307) suggest that the  
345 development of harmonic or disharmonic folds is dependent on the spacing of the competent  
346 beds, and suggest that “adjacent layers must be closer than the sum of their dominant  
347 wavelengths (perhaps as little as 10% of this distance) before they will buckle harmonically”.  
348 Numerical modelling of linear viscous multilayer folding by Schmalholtz & Mancktelow  
349 (2016, p. 1436) also indicates that more closely spaced layering may display ‘contact strain’  
350 and will tend to fold harmonically. An inspection of Fig. 3a reveals that the more competent  
351 detrital-rich markers are indeed typically closely spaced relative to fold wavelength, and  
352 therefore supports this general model.

353 The observation that individual fold packages die out both up and down their axial  
354 surface (with associated reductions in spacing of neighbouring axial surfaces) suggests that  
355 folds nucleated in the central area of each fold package and then propagated both upwards  
356 and downwards (Fig. 3a). Packages of en echelon folding generally transfer upwards through  
357 the slumped unit, with ‘jumps’ in each en echelon fold package typically occurring across  
358 thicker aragonite-rich units (Fig. 3a). Neighbouring en echelon systems tend to transfer across  
359 the same stratigraphic level marked by the weaker aragonite units (Fig. 3a), suggesting that  
360 they may form easy-slip horizons and thereby control the position and geometry of fold  
361 trains. Although en echelon segments are broadly parallel, there is a general reduction in axial  
362 surface dip together with interlimb angles up through the deformed sequence, such that the  
363 uppermost fold packages comprise sub-horizontal tight-isoclinal folds (Fig. 3a). This may  
364 suggest increasing shear towards the sediment surface during downslope-directed progressive  
365 deformation (Alsop & Marco 2013). The larger fold trains are preserved in the centre of the  
366 slump, perhaps indicating that deformation initiated here, and then transferred downwards  
367 towards the underlying basal detachment. Such patterns are similar to that encountered within  
368 thrusts of the Lisan Fm., where larger displacement may occur on thrust ramps that initiated  
369 above the basal detachment (Alsop et al. 2017a).

370

371 *Are re-fold patterns formed in slumps similar to those in metamorphic rocks?*

372 Within orogenic belts, re-fold patterns were originally considered to be created by punctuated  
373 episodes of overprinting deformation. Each deformation ‘phase’ was considered to have a  
374 distinct style and orientation that were correlated over large distances within metamorphic  
375 rocks, and resulted in re-fold patterns where fold phases interfered with one another (see  
376 discussion in Fossen 2016, p.456; Fossen et al. 2019). The strict adherence to such ‘D-  
377 number’ schemes started to break down with the realisation that heterogeneous strain could  
378 produce variable styles and orientations of structures during single progressive deformation  
379 events associated with orogenesis (e.g. Coward & Potts 1983, Holdsworth 1990; Alsop &  
380 Holdsworth 1993).

381 A similar range of re-fold geometries may be produced within slumps as observed in  
382 metamorphic rocks (e.g. Tobisch 1984; Farrell & Eaton 1987), although refolds within  
383 slumps form within a single slump event created by slope failure (Fig. 4a-j). Refolding that  
384 produces coaxial interference marked by Type 3 ‘hook’ patterns may be associated with  
385 consistent directions of flow, and a late phase of pure shear shortening that was superimposed  
386 as the slump motion reduced or stopped (Farrell & Eaton 1987). Alternatively, non-coaxial  
387 refolding marked by Type 1 ‘dome and basin’ or Type 2 ‘mushroom’ refolds implies either  
388 that the flow direction is not constant within the slump, or differential shear has locally  
389 developed within a broadly uniform slump direction leading to the creation and re-working of  
390 variably oriented folds within the flow perturbation (e.g. Alsop & Holdsworth 1993). There is  
391 no evidence within the case study that refolds are created during later loading of the slump by  
392 overburden, as the refolds are themselves truncated by erosive surfaces associated with  
393 deposition of the overlying sedimentary cap (Figs. 4j, 10). Some reworking of slump folds  
394 linked to ‘relaxation’ and reorganisation of the MTD may locally develop and have been  
395 described by Alsop & Marco (2011, p.449). There is also some evidence that refolding of  
396 slump folds developed within < 10 cm of the sediment surface may be created by the  
397 movement of the overlying water column as it ‘sloshes’ back and forth across the narrow  
398 basin during seiche events (Alsop & Marco 2012b). However, overall energy within the  
399 gravity-driven MTD will generally dissipate after the initial seismic trigger and the resulting  
400 downslope movement. In summary, the observation that a similar range of re-fold styles may  
401 be created by slump folds during SSD, compared to metamorphic folds formed during HRD,  
402 demonstrates that progressive deformation is a valid mechanism to generate such interference  
403 patterns irrespective of setting.

404

405 *What range of fold styles are created during slumping?*

406 This study has shown that all potential fold styles ranging from Class 1A, 1B, 1C, Class 2  
407 and Class 3 may form within slump systems (Figs. 5a,i, 6a-h). The development of more  
408 parallel folds (Class 1B) in competent layers, and Class 3 folds in adjacent weaker layers (e.g.  
409 Fig. 6g,h) may combine to create overall similar fold packages that “can extend for a  
410 significant distance in the profile section, in a direction parallel to the axial plane” (Price &  
411 Cosgrove 1990, p.316). Different alternating fold styles in the bilaminar (e.g. Fig. 6g, h)  
412 may therefore encourage harmonic folding. Ortner & Kilian (2016) also recognised fold  
413 styles that vary between Class 1B and Class 3 in adjacent layers within unlithified carbonates  
414 undergoing downslope-directed creep. Ramsay (1974) realised that if competent layers were  
415 thicker than the average in the multilayer, then accommodation structures must develop

416 around this thicker layer, with adjacent layers displaying ‘keel-like’ hinges (Price &  
417 Cosgrove 1990, p.320). Similar structures are observed in fold hinges formed next to thicker  
418 detrital layers in the present study (e.g. Fig. 8p).

419 Waldron & Gagnon (2011) suggested that the style and class of folds in different  
420 lithological layers may be critical in identifying SSD. Unlithified muds may display buckle  
421 folds (Class 1B), while adjacent sands are weaker (due to greater porosity and water content)  
422 and may display Class 2 (similar) folds. These relationships in unlithified sediments are the  
423 opposite to that typically encountered in lithified rocks, where sandstones are normally more  
424 competent than mudstones during folding associated with HRD. Indeed, Hibbard & Karig  
425 (1987, p.848) record Class 1B folds in sandstones and Class 2 or Class 3 folds in mudstones  
426 and suggest that beds may have been lithified at the time of folding in a variably deformed  
427 accretionary complex in SW Japan. In summary, although folding in unlithified sediments is  
428 formed via deformation associated with ‘hydroplastic’ particulate flow, rather than  
429 recrystallization as in metamorphic hard rock deformation (HRD), the range of geometries  
430 produced are similar to one another. This includes the development of axial planar fabrics in  
431 slump folds as observed previously within the Lisan Fm. (e.g. Alsop & Marco 2014; Alsop et  
432 al. 2019a, b). Different lithologies may however behave more or less competently during  
433 HRD or SSD and thereby create folds of distinct styles in each setting.

434

435 *How are slump folds subsequently modified?*

436 Post-buckle flattening

437 Ramsay suggested (1967, p.411) that Class 1B parallel folds may be subsequently modified  
438 by ‘flattening’ that resulted in Class 1C geometries. Ramsay (1967, p.433) went on to suggest  
439 that “Many of the apparently similar folds seen in naturally deformed rocks probably owe  
440 their initial development to the buckling mechanism”. We have therefore undertaken analysis  
441 of such post-buckle modification in order to test if this is an appropriate mechanism to create  
442 the range of fold styles within the case study. Homogeneous flattening normal to the axial  
443 plane of an original Class 1B fold results in an increased fold amplitude and a Class 1C fold  
444 (e.g. Farrell & Eaton 1987, p.193), whereas flattening parallel to the axial plane leads to a  
445 reduction in fold amplitude and a Class 1A fold (e.g. Twiss & Moores 2007, p.374).  
446 Hudleston (1973) recognised that Class 1C folds may be produced during simultaneous  
447 buckling and flattening, and there is therefore no necessity for separate and sequential  
448 ‘phases’ of fold development, with folds created during progressive deformation within the  
449 slump.

450 Our analysis reveals that if flattening were the sole cause of the different fold styles,  
451 then it would need to form a significant element in the deformation. Using the technique of  
452 Lisle (1992) described previously, elliptical ratios of ~1.7 are calculated, suggesting a relative  
453 vertical shortening of 30% (Fig. 7a-d). In other cases, some Class 1C folds defined by detrital-  
454 rich markers would in fact suggest up to ~50% homogenous shortening (Fig. 7i-l). The ‘mis-  
455 match’ of axial surfaces that ‘jump’ when traced across weak beds may be a consequence of  
456 such flattening (e.g. Fig. 3a). However, recumbent crenulations observed in aragonite laminae

457 in the hinges of upright Class 1A folds record only limited vertical shortening (see Alsop &  
458 Marco 2011). Such horizontal crenulations were attributed to ‘compaction’ by Alsop &  
459 Marco (2011). As there is no other evidence for significant flattening, with thrusts preserving  
460 pristine dip angles of 30° (e.g. Alsop et al. 2017a, b), and a complete absence of folded  
461 clastic dykes (Alsop et al. 2019a), this amount of vertical shortening seems unfeasible. Post-  
462 buckle flattening may represent a small component that affects upright Class 1B folds close  
463 to the sediment surface where uncompacted sediment was most water-saturated, and results  
464 in locally thickened limbs and thinned hinges (Class 1A) (Fig. 10). However, given the lack  
465 of supporting evidence, post-buckle flattening should not perhaps be applied to the entire  
466 sequence and we therefore suggest that other mechanisms may also have operated to modify  
467 fold shapes.

468

#### 469 Post-buckle shearing

470 The relationship between layer thickness and fold wavelength suggests buckling is the  
471 dominant fold mechanism within slump folds of this case study and elsewhere (e.g.  
472 Woodcock 1976b; Farrell & Eaton 1987). As noted above, deviation of buckle folds from  
473 Class 1B may be a consequence of vertical flattening, creating upright Class 1A with  
474 thickened limbs (e.g. Figs. 6a,b, 10), or recumbent Class 1C folds with thinned limbs (e.g.  
475 Fig. 6c,d) (e.g. Farrell & Eaton 1987, their figure 11). However, the variation in geometry  
476 may also be created by simple shear being imposed on the buckles resulting in thickened or  
477 thinned fold limbs depending on the orientation of the fold and amount of shear (e.g. Farrell  
478 & Eaton 1987; see discussion in Twiss & Moores 2007).

479 It has previously been suggested that buckle folds within the Lisan Fm. may form  
480 during layer-parallel shortening associated with pure shear, and that fold initiation may also  
481 potentially involve density-driven flow into antiforms associated with Rayleigh-Taylor  
482 instabilities (e.g. Alsop & Marco 2011, p.453) (Fig. 10). Downslope-directed simple shear is  
483 then progressively imposed on these folds to create the observed range of geometries (Alsop  
484 & Marco 2011 p.446) (Fig. 10). The tightening and rotation of upright buckles to overturned  
485 folds in the case study (Fig. 5e,f,g) is consistent with increasing downslope-directed non-  
486 coaxial dominated deformation (e.g. Woodcock 1976b; Farrell & Eaton 1987; Twiss &  
487 Moores 2007, p.360; Dasgupta 2008) (Fig. 10). This interpretation is also supported by  
488 numerical simulations of downslope-directed gravity sliding that results in buckle folds with  
489 hinges that rotate towards the flow direction and axial planes that sequentially rotate into the  
490 plane of flow during continued slumping (Schmalholz & Schmid 2012, p. 1817).

491 Recumbent folds are sometimes marked by thicker overturned lower limbs compared  
492 to the upper fold limbs (e.g. Fig. 3a, 5h,i). Folds with thicker lower limbs suggest that  
493 original buckle was upright, and that the downslope limb has spent some time in the  
494 contractional field of the strain ellipsoid (while the upslope limb was always in the  
495 extensional field and therefore thinner). In other cases, the lower fold limbs are more thinned  
496 (Fig. 6g,h), suggesting that these buckles may have originated with a downslope asymmetry  
497 with shorter, steeper limbs marking vergence towards the basin (see Alsop & Marco 2013,

498 their fig. 7). The observation that most folds in the case study display thinned lower limbs is  
499 consistent with original buckles displaying a downslope-directed asymmetry and vergence  
500 towards the basin. In addition, we do not generally observe recumbent Class 1B or upright  
501 Class 1C or Class 2 folds, suggesting that folds are modified as they become overturned  
502 during downslope-directed shearing. These observations collectively suggest that slump folds  
503 originated as more upright buckles that have been subsequently deformed by downslope-  
504 directed simple shear deformation (see Farrell & Eaton 1987; Alsop & Marco 2011) (Fig.  
505 10). There may therefore be a combination of post-buckle flattening and shearing, with  
506 recumbent Class 1C being created by flexural flow rather than flattening of original Class 1B  
507 buckles.

508

#### 509 Post-buckle flow

510 Price & Cosgrove (1990, p.402) note that “post-buckle flattening may involve physical  
511 migration of the relatively incompetent material”. They further note (Price & Cosgrove 1990,  
512 p.403) that incompetent layers in the short limbs of folds may have thinned due to “the  
513 migration of material from the limbs to the hinge region in response to stress gradients, which  
514 tend to become greater during the last stages of fold amplification”. Although our analysis is  
515 performed on unlithified detrital-rich beds, these beds are exceptionally weak and capable of  
516 flow when compared to lithified rocks.

517 Within the case study slump folds, some upright synforms are markedly thickened in  
518 the hinges when compared to adjacent antiforms (Fig.7e, 8m), while the converse is rarely  
519 observed. This suggests that in some cases there may be a component of material flow from  
520 the antiformal crest down into the synformal troughs. The typical position of Class 1A folds  
521 towards the upper parts of slumps directly below erosive truncation surfaces linked to the  
522 sediment water interface, suggests that gravity-driven migration of sediment within steep  
523 water-saturated fold limbs close to the surface may be a significant factor (e.g. Figs. 7, 10).  
524 Farrell & Eaton (1987, p.186) examined slump folds in Cyprus and the Spanish Pyrenees and  
525 also note “flowage of sediment into cores of folds”, which on examination of their figure 3d  
526 is predominantly towards the troughs of synforms. Flow of sediment down off earlier  
527 structural highs and folds was previously termed ‘relaxation’ by Alsop & Marco (2011), and  
528 may form a significant local factor in fold modification.

529

#### 530 Summary

531 A range of post-buckle processes may have contributed to the final geometry of the slump  
532 folds. Thickening of upright fold limbs to create Class 1A folds suggests that a component of  
533 post-buckle vertical ‘flattening’ has operated during the slump process, while thickening of  
534 synformal hinges relative to neighbouring antiforms is consistent with gravity-driven  
535 hydroplastic flow of sediment from crests to troughs (Fig. 10). Although buckle folds may  
536 have initiated during layer-parallel shortening associated with pure shear, they were  
537 subsequently modified by downslope-directed simple shear resulting in overturned Class 2  
538 folds (Farrell & Eaton 1987; Alsop & Marco 2011, 2013) (Fig. 10). In addition, other more

539 nebulous influences such as locally increasing fluid pressure during the folding process may  
540 lead to a relative weakening of competent layers (e.g. Price & Cosgrove 1990, p.295),  
541 thereby encouraging more Class 1C and Class 2 similar folds to develop during the evolution  
542 of the slump. Overall volume loss via expulsion of fluids (e.g. Price & Cosgrove 1990) and  
543 increased grain packing may also influence the geometry of folds, although is difficult to  
544 quantify. These processes may collectively modify the initial fold shape to create the range of  
545 fold geometries that are currently observed in the study area.

546

547 *Are estimates of viscosity contrasts in slumps folds similar to those in metamorphic rocks?*

548 Schmalholz & Podladchikov (2001) have demonstrated via strain contour maps that viscosity  
549 contrasts between layers of between 50 and 250 are typical in metamorphic rocks  
550 (Schmalholz & Podladchikov 2001; see also Hudleston & Treagus 2010). These values are  
551 however observed to vary depending on lithological controls (e.g. Druguet et al. 2009, their  
552 fig. 3). Schmalholz & Podladchikov (2001, p. 211) also note that the strain contour map is  
553 applicable not just to buckles created during pure shear, but also to asymmetric folds that  
554 formed during simple shear. In general, the values of viscosity contrast calculated from the  
555 slump folds in the study area display a similar range to that obtained from metamorphic rocks  
556 (Fig. 8, 9).

557 Thus, similar geometries between metamorphic folds created during HRD and slump  
558 folds formed during SSD are borne out by a similar range of viscosity contrasts on strain  
559 contour maps. However, folds within the case study slumps form before the overlying  
560 sedimentary cap was deposited out of suspension “in a matter of just hours or days” (Alsop et  
561 al. 2016, p.80) and absolute viscosities in slumps will therefore be much lower. As a guide,  
562 samples of modern Dead Sea sediment yield viscosity measurements of just 0.3 Pa.s and 3  
563 Pa.s at depths of 10 cm and 30 cm respectively below the present sediment surface (Wetzler  
564 et al. 2010, p.304), while Levi et al. (2008) suggest that detrital layers in the Lisan Fm. have  
565 dynamic viscosity values in the range of 0.03–0.3 Pa.s. In terms of strain rates, Price &  
566 Cosgrove (1990, p.369) estimated the time taken during HRD to create folds of 500 m  
567 wavelength, based on typical plate tectonic velocities of ~3.3cm/year, and suggested strain  
568 rates of  $10^{-13}$ /s. If we make a crude assumption within the case study that a 1 m wavelength  
569 slump fold is formed in 1 day, (Alsop et al. 2016), then suggested strain rates during SSD are  
570 of the order of  $10^{-5}$ /s, while a similar fold forming more slowly over 1 week gives strain rates  
571 of  $10^{-6}$ /s. The conclusion is that, despite the geometrical similarity and similar range of  
572 viscosity contrasts between metamorphic folds and SSD folds, the strain rates within slumps  
573 that create MTDs are 7-8 orders of magnitude greater and may be viewed as geologically  
574 instantaneous.

575 There are a number of considerations when estimating viscosity contrasts and total  
576 shortening within beds defining slump folds. Price & Cosgrove (1990, p.299) note that the  
577 amount of bulk layer-parallel shortening that develops before buckles start to amplify  
578 increases as the viscosity contrast decreases. This idea was developed further by Treagus  
579 (1997) who suggested that there may be between 5 – 20% bulk (layer-parallel) shortening

580 before buckles start to amplify, with this pre-buckle shortening most significant in modelled  
581 folds with lower viscosity ratios (Treagus 1997; Hudleston & Treagus 2010, p.2059). Teixell  
582 & Koyi (2003) and Koyi et al. (2004) also note that ~20% layer-parallel shortening may  
583 develop in models undergoing folding and thrusting. The net result may be that Schmalholz  
584 & Podladchikov (2001) strain plots underestimate total shortening as they fail to take account  
585 of this layer-parallel component.

586 It has also been suggested that for a regular fold wave-train with a dominant wavelength to  
587 develop, the viscosity-ratio should be >100:1, otherwise layer-parallel shortening (noted  
588 above) that pre-dates buckling becomes so large that it masks and would 'overshadow  
589 folding' (Biot 1957, see discussion in Price & Cosgrove 1990, p. 281). However, our results  
590 suggest that in some cases, fold wave-trains may display viscosity ratios of significantly less  
591 than 100:1 (e.g. Fig. 8c), and in some instances less than 50:1 (Fig. 8o). It has been recognised  
592 from experiments that with increased homogeneous layer-parallel shortening that  
593 accompanies the buckling, then wave-trains with dominant wavelengths may indeed form at  
594 significantly lower viscosity contrasts of 25:1 (Hudleston 1973). Our results are therefore  
595 consistent with a component of homogeneous layer shortening that accompanies the buckling  
596 of competent layers during the slump process. Further details of the role of viscosity ratios in  
597 the folding process are provided by Schmid et al. (2010) and more recently are reviewed by  
598 Schmalholz & Mancktelow (2016).

599 The correlation we describe between fold wavelength ( $\lambda$ ) and amplitude ( $A$ ) (e.g. Fig.  
600 9a-c) is in general agreement with Wetzler et al. (2010) who correlated  $A^2$  and wavenumber  
601 (i.e.  $1/\lambda$ ) in slump folds from the Lisan Fm. Wetzler et al. (2010) suggest that the observed  
602 folds form part of a continuum of structures that initiate as linear waves during Kelvin-  
603 Helmholtz instability and evolve in to folds, finally culminating in breccias associated with  
604 turbulence and instability (also see Alsop & Marco 2011, p.453). Analysis of fold wavelength  
605 ( $\lambda$ ) to layer thickness ( $h$ ) from the measured slump folds (e.g. Fig. 8a) suggests a general ratio  
606 of ~5:1 (Fig 9d,e,f). This is significantly less than the 27:1 ratio calculated by Twiss & Moore  
607 (2007, p.562) for folds in lithified rocks. Although this  $\lambda/h$  value is viewed as a constant, it is  
608 recognised that wavelength is not a very sensitive function, as slight changes result in  
609 significant differences to the estimate of the viscosity contrast. In addition, if there is a  
610 component of multilayer folding (i.e. layers are more closely spaced) then the wavelength  
611 produced will be less than for an equivalent thickness of a single layer buckle fold (Price &  
612 Cosgrove 1990, p. 310). The net effect of this multilayer folding would be to increase both  
613 the  $A/\lambda$  and  $h/\lambda$  values. This will result in an overestimate of the % contraction, and a  
614 potential over- or under estimate of the viscosity contrast on the Schmalholz & Podladchikov  
615 (2001) strain maps designed for single layer folds. The role of multilayer spacing on the  $\lambda$  to  
616  $h$  ratios of individual folded layers have been recently summarised by Schmalholz &  
617 Mancktelow (2016, p.1437).

618 It is interesting to note that buckle fold layers that are entirely isolated within  
619 aragonite-rich units may display greater viscosity contrast as % contraction increases (e.g.  
620 Fig. 8a-c), whereas if layers are closer together, there is an apparent reduction in viscosity  
621 contrast with increasing % contraction resulting in less steep trends on strain contour plots  
622 (Fig. 8f, i, l, o). This apparent discrepancy may reflect a number of variables. If a measured  
623 single layer fold actually forms part of a multi-layer, the multilayer will act as a thicker

624 mechanical unit, with a greater effective thickness, meaning that measured  $h/\lambda$  values would  
625 be under estimated. This multilayer influence may be most significant where folds initiate  
626 and amplitudes are lower (with lower % contraction) and may result in some data sets  
627 defining less steep trends on strain plots i.e. the effect of multi-layers may be to overestimate  
628 viscosity contrasts at lower values of % shortening (see also Druguet et al. 2009, p.503). An  
629 alternative interpretation would be that viscosity contrasts actually vary as folds develop due  
630 to expulsion of fluids during folding. Price & Cosgrove (1990, p.398) examined folds formed  
631 during HRD and argued that “the fold acts as a pump which, if the external energy is  
632 sufficient, continues to extrude fluids from itself until it is almost dry”. Any expulsion of  
633 fluids leading to variation in viscosity contrasts would have the greatest effect on resulting  
634 fold geometries if it occurred relatively early in the folding process when amplitudes are  
635 smaller and before fold wavelengths are established in the system. We here suggest that  
636 during SSD, a preferential expulsion of fluids from weaker (aragonite-rich) layers as the %  
637 shortening increases, would lead to an increase in viscosity of these layers and consequently a  
638 reduction in the contrast between viscosities of the aragonite and detrital-rich layers.

639 In summary, the overall correlation between bed thickness ( $h$ ), fold wavelength ( $\lambda$ )  
640 and amplitude ( $A$ ) suggests that buckling is the dominant fold mechanism in the slump folds  
641 of the study area (Fig. 8, 9). Strain maps developed by Schmalholz & Podladchikov (2001)  
642 are equally applicable to folds created during HRD of metamorphic rocks and slump folds  
643 formed during SSD, and suggest similar viscosity contrasts ranging from 50-250. Within  
644 slump folds, components of potential layer-parallel shortening would lead to an under-  
645 estimate of total shortening, while an element of multi-layer folding may lead to an  
646 overestimate of the % contraction and less accurate estimates of viscosity contrasts. We  
647 speculate that it may also be possible for viscosity contrasts to actually vary as fold trains  
648 evolve due to preferential expulsion of fluids as folds tighten.

649

## 650 **Conclusions**

651 In this study, classical techniques of structural analysis developed by John Ramsay amongst  
652 others during the investigation of folds formed in metamorphic terranes have been applied to  
653 folds generated within unlithified soft-sediments. Using superbly exposed slump folds  
654 created during SSD within MTDs of the Lisan Fm., we have demonstrated that recumbent  
655 fold packages define en echelon fold trains that display a harmonic style of folding. Slump  
656 folds are entirely systematic and should not be viewed as ‘chaotic’, although upright buckles  
657 may form disharmonic folds that cannot be traced for significant distances along their axial  
658 surfaces before truncation by overlying sedimentary caps. We also recognise a similar range  
659 of refold patterns created during SSD within slump folds when compared to HRD in  
660 metamorphic rocks. Refolds may also be truncated by the overlying sedimentary cap,  
661 indicating that they formed during a single progressive slump event and are not a product of  
662 superposition of later ‘events’ (Fig. 10).

663 A wide variety of fold styles ranging from Class 1A, 1B, 1C, Class 2 and Class 3  
664 geometries are developed within slumps of the case study. Within individual folds, aragonite-  
665 rich layers may display Class 2 (similar) folds, while adjacent detrital-rich beds are marked  
666 by Class 1C fold geometries. Such varying styles of folding allow individual folds to



667 propagate further along their axial planes without encountering significant accommodation  
668 problems around the fold hinge. Upright buckle folds defined by detrital-rich marker beds  
669 display a Class 1B (parallel) style of folding, while recumbent folds display Class 2 (similar)  
670 fold styles. Upright fold hinges and axial planes may undergo rotation during progressive  
671 simple-shear dominated deformation as the MTD moves downslope under the influence of  
672 gravity. In detail, thickening of upright fold limbs relative to the hinge to form Class 1C folds  
673 suggests a component of vertical flattening that post-dates original folding, but pre-dates  
674 deposition of the overlying sedimentary cap shortly after slope failure. This flattening appears  
675 to be most significant towards the top of each MTD, where sediments are more water-  
676 saturated (Fig. 10). In addition, upright folds are also subsequently modified by flow of  
677 sediment from antiformal crests to troughs of synforms, together with simple shear  
678 deformation associated with continued downslope movement that results in Class 1C and  
679 Class 2 similar folds displaying recumbent attitudes (Fig. 10).

680 Our study has involved the measurement of fold wavelength ( $\lambda$ ), amplitude ( $A$ ) and  
681 bed thickness ( $h$ ) in order to estimate % shortening and viscosity contrast between folded  
682 layers on strain contour maps (Schmalholz and Podladchikov, 2001). Estimates of viscosity  
683 contrast from folded layers within metamorphic rocks formed during HRD and slump folds  
684 created during SSD both suggest broadly similar values in the range of 50-250. We speculate  
685 that it may also be possible for viscosity contrasts to vary as fold trains evolve due to  
686 preferential expulsion of fluids as folds tighten. The geologically instantaneous nature of  
687 slope failure, which may be 7-8 orders of magnitude greater than folding associated with  
688 HRD, indicates that the absolute viscosities will however be considerably lower in slumps.

689

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698

## 699 **Figure Captions**

700 **Fig 1.** a) Tectonic plates in the Middle East. General tectonic map showing the location of the  
701 present Dead Sea Fault (DSF). The Dead Sea Fault transfers the opening motion in the Red  
702 Sea to the Taurus-Zagros collision zone. b) Generalised map (based on Sneh and Weinberger  
703 2014) showing the current Dead Sea including the position of the Masada, Peratzim and Zin  
704 localities referred to in the text. The extent of the Lisan Fm. outcrops are also shown, together  
705 with the general slump directions of the MTD's around the basin.

706 **Fig 2** a, b) Photographs of light-coloured aragonite-rich laminae and dark detrital-rich  
707 laminae forming the Lisan Fm. An infilling sedimentary cap that blankets underlying  
708 structures at Zin is shown in a), while b) is a close-up image of undeformed laminae from  
709 Peratzim. Refer to Fig. 1b for details of locations. Pairs of photographs (c, e, g) and  
710 associated stereonet (d, f, h) of NE-verging fold and thrusts developed within slump 4 at  
711 Peratzim (see Alsop et al. 2016). In stereonet, fold hinges (solid red circles), poles to axial  
712 planes (open blue squares), and associated thrusts are shown as red great circles and poles as  
713 solid red squares. Calculated slump transport is towards the NE (see Alsop & Marco 2012a).  
714 All photographs marginally overlap with one another and have NE on the right-hand side for  
715 consistency.

716 **Fig. 3** Harmonic and disharmonic fold analysis from folds in Masada (see Fig. 1b for  
717 location). a) Recumbent en echelon folds verging towards the East. The half wavelength of  
718 folds is approximated by the spacing (S) of adjacent axial surfaces (shown in red and blue),  
719 while the continuity of the axial surface may be directly measured (D), to produce a ratio (H)  
720 which is equivalent to D/S (see Twiss & Moores 2007, p.290). b) Stereonet of folds from  
721 Masada shown in a) with transport direction towards  $074^\circ$  and key to stereonet given below.  
722 c) Graph showing D plotted against S. d) Graph showing H plotted against S. e) Photograph  
723 of individual Masada locality with stereonet f), and plots of H, D and S parameters (g, h). i)  
724 Photograph of individual Masada locality with stereonet j), and plots of parameters (k, l). In  
725 all stereonet, fold hinges (solid red circles), mean fold hinge (open red circle), poles to axial  
726 planes (solid blue squares), mean axial plane (blue great circle and pole as open blue circle).  
727 Great circle and pole to sub-horizontal bedding are shown in green in f, j).

728 **Fig. 4** Photographs (a-d) of minor folds within Moine psammities from the classic locality by  
729 Loch Monar Dam in NW Scotland [UK Grid NH1989 3882] used by Ramsay (1962, 1967) to  
730 establish a refold classification scheme. a) Plan view of intense minor folding within  
731 interlayered pelites and psammities. Adjacent axial traces (highlighted in blue and red)  
732 define harmonic folds. b, c) Plan view of Type 1 dome and basin fold interference patterns  
733 with bedding in psammite (highlighted in yellow) defining closed outcrop patterns. d) Plan  
734 view of Type 3 hook fold interference patterns with axial traces of older folds (highlighted in  
735 red) being refolded by younger folds (highlighted in blue). Photographs (e-j) of sections  
736 through refolds within slumps at Peratzim (See Fig. 1b for location). e, f) Type 1 dome and  
737 basin refold pattern, with coloured pencils parallel to the fold hinges in e). g) Type 2  
738 'mushroom' refold pattern. h, i, j) Type 3 hook interference patterns. In j), the refold is  
739 truncated by an overlying erosive surface and sedimentary cap. 10 cm long chequered rule  
740 and 15 mm diameter coin for scale.

741 **Fig 5** a) Dip isogons join points of equal dip on adjacent folded surfaces,  $t_0$  is layer thickness  
742 measured along the axial surface, while  $t_\alpha$  is orthogonal layer thickness measured at various  
743 angles ( $\alpha$ ) to the reference plane. b) Class 1 folds are marked by convergent dip isogons, c)  
744 Class 2 folds by parallel dip isogons, and d) Class 3 folds by diverging dip isogons. e) Single-  
745 layer train of buckle folds defined by a detrital-rich bed at Peratzim (see Fig. 1b for location).  
746 Buckle folds display a progressive tightening coupled with a rotation and reduction in dip of  
747 the axial surfaces towards the downslope direction. f) Excavation of buckle folds shown in e)  
748 reveals 3D variation in fold hinge orientations marked by different coloured pencils. 10 cm

749 chequered rule for scale. g) Stereonet of upright folds (in red) and more overturned folds with  
 750 inverted lower limbs (in blue) shown in e) and f). Axial planes are shown as great circles, and  
 751 poles as solid squares. Fold hinges are shown by solid circles. h) Close-up photograph of  
 752 buckle folds that display thickened overturned limbs relative to the upper limbs (see Fig. 5e  
 753 for position). Representative 70° and 45° dip isogons are drawn on the upper and lower limbs  
 754 of the fold. i)  $t'_\alpha$  graph (where  $t'_\alpha = t_\alpha / t_0$ ) plotted against dip angle ( $\alpha$ ) to create a series of  
 755 fold classes (Ramsay 1967, p. 366).  $t'$  alpha graph relates to dip-isogon analysis of detrital-  
 756 rich (blue) and aragonite-rich (red) layers around buckle fold shown in h). Data is divided  
 757 into upper fold limbs (squares) and lower fold limbs (circles).

758 **Fig. 6** Paired photographs and  $t'$  alpha ( $\alpha$ ) plots relating to dip-isogon analysis of different  
 759 layers around folds (a-b, c-d, e-f, g-h). Dip isogons join points of equal dip on adjacent folded  
 760 surfaces,  $t_0$  is layer thickness measured along the axial surface, while  $t_\alpha$  is orthogonal layer  
 761 thickness measured at various angles ( $\alpha$ ) to the reference plane (see Fig. 5a). Upper fold  
 762 limbs are represented by coloured circles and lower limbs by squares. Photographs a) and c)  
 763 are from Peratzim, while e) and g) are from Zin.

764 **Fig. 7** Analysis of post-buckle flattening in folds developed directly beneath sedimentary  
 765 caps. In each case, paired photographs and associated inverse layer thickness analysis based  
 766 on Lisle (1992) are shown in a-b) from Peratzim, e-f) from Masada and i-j) from Peratzim. a)  
 767 and b) illustrate the inverse thickness method, which plots  $(1/t)$  for various orientations of the  
 768 layer tangent around the fold, where  $t$  is the orthogonal thickness between tangents for the  
 769 inner and outer layer boundaries (see Lisle 1992, p.370). The inverse thickness  $(1/t)$  is plotted  
 770 from a common central point, each in the direction of the tangent to create an array of points  
 771 to which a best-fit ellipse is matched. Examples of upslope (SW in red) tangents drawn at  
 772 70°, and downslope (NE in blue) tangents drawn at 45° are shown in a) and b). Photographs  
 773 and ellipses are 'unstrained' back to circles using the method of Srivastava and Shah (2006),  
 774 and are shown in the right-hand side in c-d), g-h) and k-l). In each case, data collected from  
 775 the western or SW (upslope) limb is shown in red, while the eastern or NE (downslope) limb  
 776 is shown in blue. The change in area from ellipse to restored circle does not represent a  
 777 change in volume of the sediment.

778 **Fig. 8** Photographs, associated graphs showing wavelength plotted against amplitude /  
 779 wavelength, and strain contour maps of Schmalholz & Podladchikov (2001) for a series of  
 780 individual fold trains (a-c, d-f, g-i, j-l, m-o, p-r). In e) and f), blue squares represent data  
 781 collected from the same bed in an adjacent outcrop of the slump. In a), wavelength ( $\lambda$ ) is  
 782 defined as the distance between two points that occupy a similar position on the fold train  
 783 (i.e. between adjacent synform hinges). Thickness of a layer ( $h$ ) is measured orthogonal to the  
 784 folded bed, while amplitude ( $A$ ) is defined as half the distance from the trough to the crest of  
 785 upright folds. Strain contour plots of Schmalholz & Podladchikov (2001) show estimated %  
 786 shortening and viscosity contrasts for folded layers. In each of the graphs, different coloured  
 787 symbols represent different individual fold trains. Arrows show general trends of data on  
 788 each graph, while circled letters correspond to individual folds marked on adjacent  
 789 photographs. Yellow notebook (21 cm long), 10 cm chequered rule, and 15 mm diameter  
 790 coin act as scales on photographs,

791 **Fig. 9** a-c) Graphs showing fold amplitude plotted against fold wavelength for a) minor folds  
 792 (< 200 mm wavelength), b) mesoscopic folds (<1000 mm wavelength) and c) overall data on  
 793 a log-log plot. d-f) Graphs showing folded layer thickness against fold wavelength for d)

794 minor folds (< 200 mm wavelength), e) mesoscopic folds (<1000 mm wavelength) and f)  
795 overall data on a log-log plot. g-h) Graphs showing folded layer thickness against fold  
796 amplitude for g) minor folds (< 100 mm amplitude), b) mesoscopic folds (<500 mm  
797 amplitude). i) Strain contour map of Schmalholz & Podladchikov (2001) showing estimated  
798 % shortening and viscosity contrasts for folded layers. In each of the graphs (a-i), different  
799 coloured symbols represent different individual fold trains, including those shown in Figure 8.

800 **Fig. 10** Schematic summary cartoon highlighting the main geometries and styles of folding  
801 observed during soft-sediment deformation of competent (brown) markers above a basal  
802 detachment. Layer-parallel shortening associated with pure shear creates upright Class 1B  
803 buckles (left of diagram) that undergo vertical flattening (large blue arrow) and flow of  
804 sediment down fold limbs towards synformal troughs to create Class 1A antiforms. A  
805 potential component of density-driven flow up into antiforms may be associated with  
806 Rayleigh-Taylor instabilities. Increasing downslope-directed simple shear deformation  
807 towards the right of the figure results in rotation of antiformal and synformal axial surfaces  
808 (red and blue dashed lines respectively) towards the flow plane, resulting in overturned and  
809 harmonic folds with Class 1C or Class 2 (similar) geometry. Refolded folds created during  
810 the progressive downslope-directed shearing are truncated by the overlying erosive  
811 sedimentary cap (orange bed) that is deposited from suspension following slope failure.

812

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